

Ground-based detection of gravitational waves from intermediate-mass-ratio inspirals

Ilya Mandel
Northwestern University

17/09/2008 @ IofA, Cambridge

based on

Mandel, Brown, Gair, Miller: ApJ 681 1431 (2008)

Gair, Li, Mandel: PRD 77 024035 (2008)

Mandel: arXiv:0707.0711

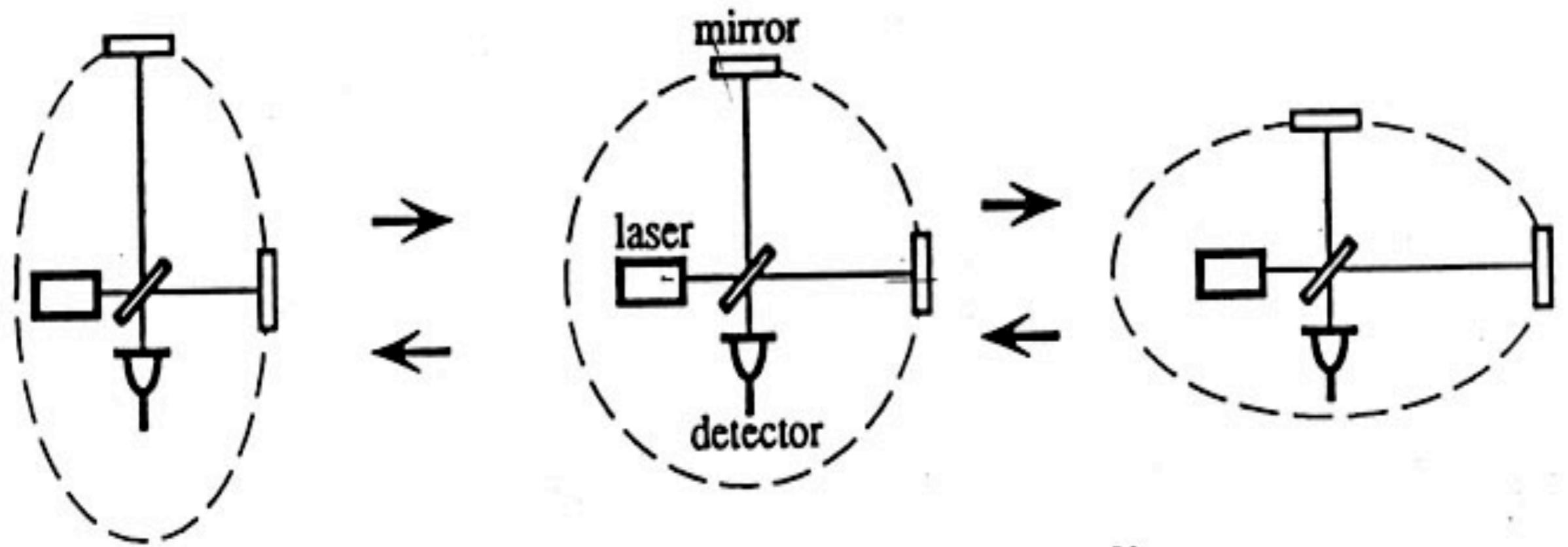
Outline

- GWs in ground-based detectors
- Intermediate-mass-ratio inspirals into intermediate-mass black holes: rates and characteristics
- Probing the strong field region near a black hole

Outline

- GWs in ground-based detectors
- Intermediate-mass-ratio inspirals into intermediate-mass black holes: rates and characteristics
- Probing the strong field region near a black hole

GW detection with Michelson-type interferometers



LIGO



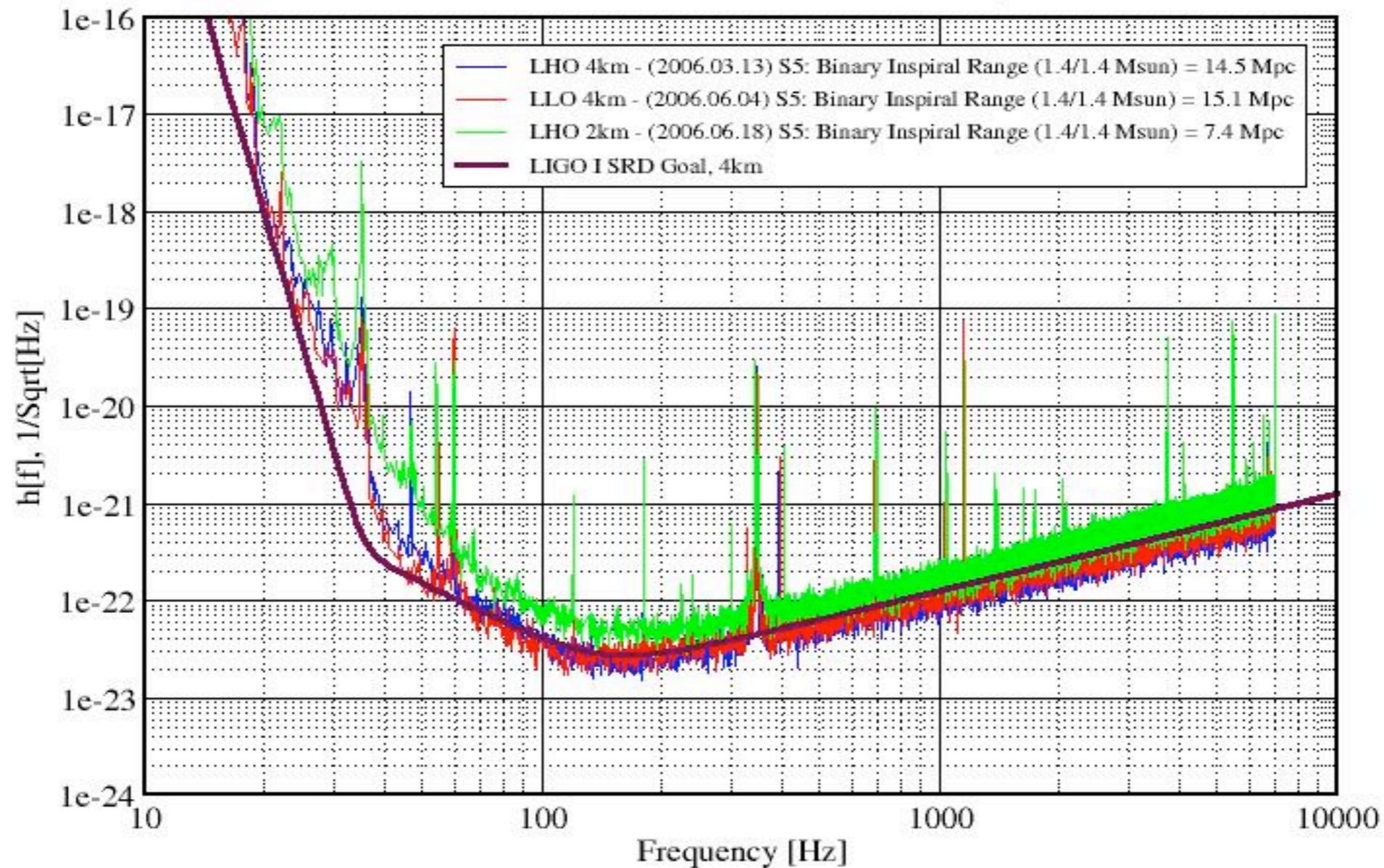
- 4 km long arms
- Typical strains $h = \Delta L / L \sim 10^{-22}$
- Needs to measure $\Delta L = hL \sim 10^{-17}$ m

LIGO Noise Curve

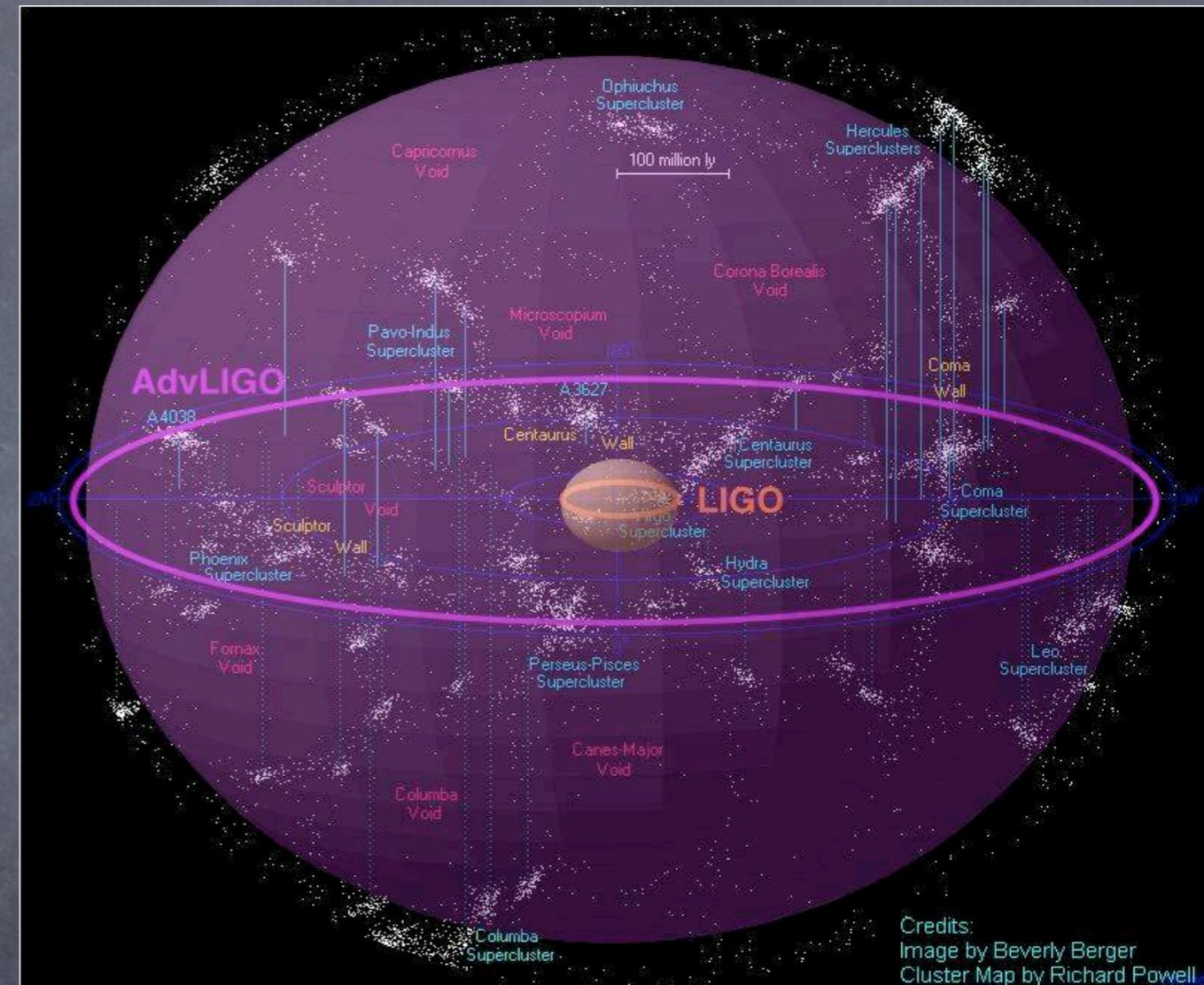
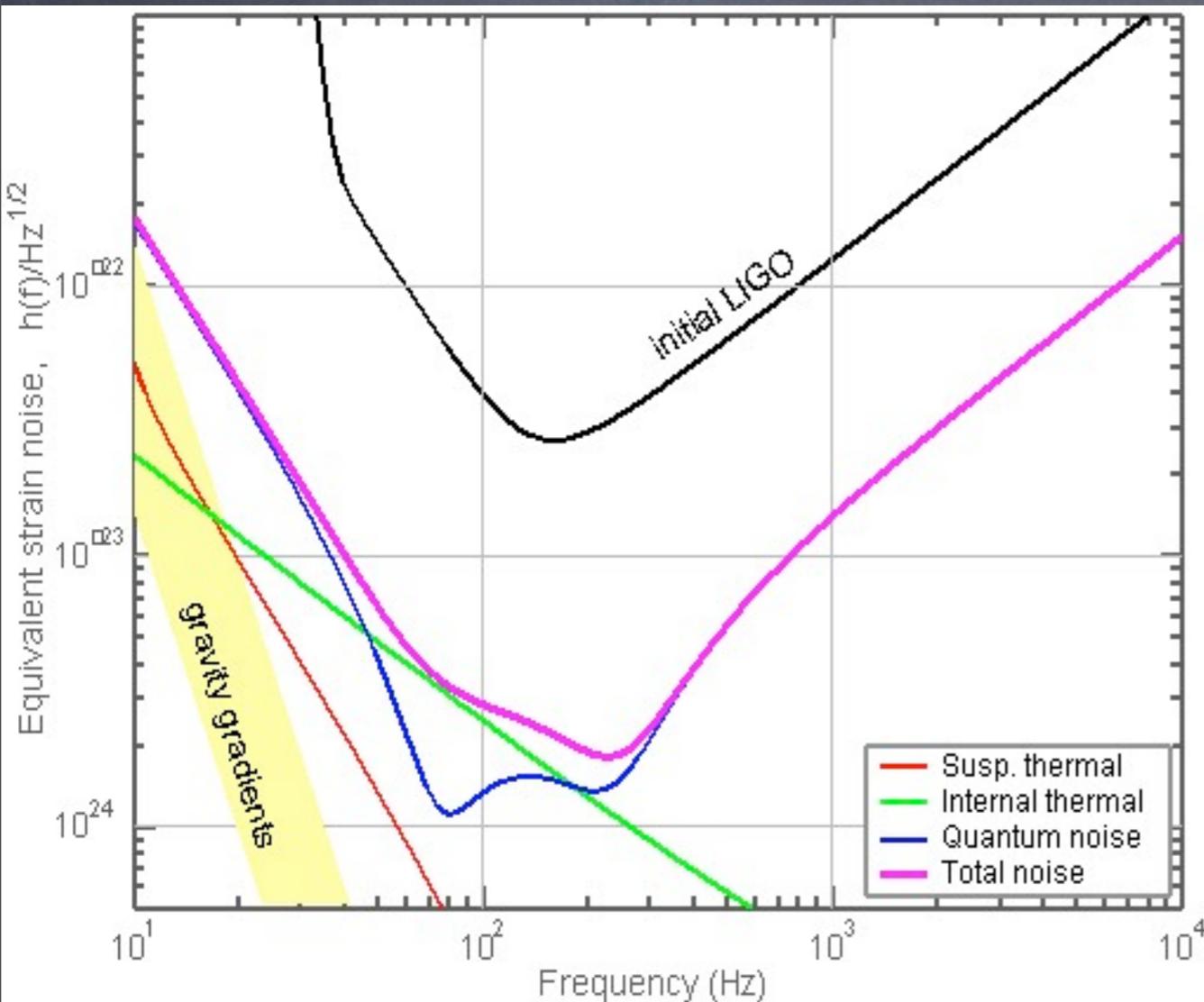
Strain Sensitivity for the LIGO 4km Interferometers

S5 Performance - June 2006

LIGO-G060293-01-Z



Advanced LIGO



- x10 in range -> x1000 in event rate
- 10 Hz low frequency cutoff

Outline

- GWs in ground-based detectors
- Intermediate-mass-ratio inspirals into intermediate-mass black holes: rates and characteristics
- Probing the strong field region near a black hole: testing the no-hair theorem

Intermediate-mass-ratio inspirals (IMRIs)

- IMRIs have mass ratios between 10 and 10^4
- LIGO IMRIs: Inspirals of compact objects (1.4 solar-mass Neutrons Stars to 10 solar-mass Black Holes) into intermediate mass black holes (IMBHs, 50–350 solar masses)
- Indirect evidence for IMBH existence in globular clusters (50 – 10^4 solar masses)
 - Observational evidence (e.g. Macarone et al.)
 - Simulations (e.g. McMillan et al., O'Leary et al.)
 - Simulations vs. Observations (e.g. Trenti)
- IMRIs could be the first proof of IMBH existence!

Event Rates: Mechanisms

- Three-body interactions: IMBH swaps into binaries, forms CO-IMBH binaries which are tightened via three-body interactions with other stars, then merge via GW radiation reaction
- Direct capture via energy loss to GWs
- Kozai resonances in hierarchical triple systems: inner binary eccentricity is driven up by outer companion
- Tidal capture of MS star that evolves into CO while in orbit
- Tidal interactions (orbital-vibrational coupling) for NS inspirals

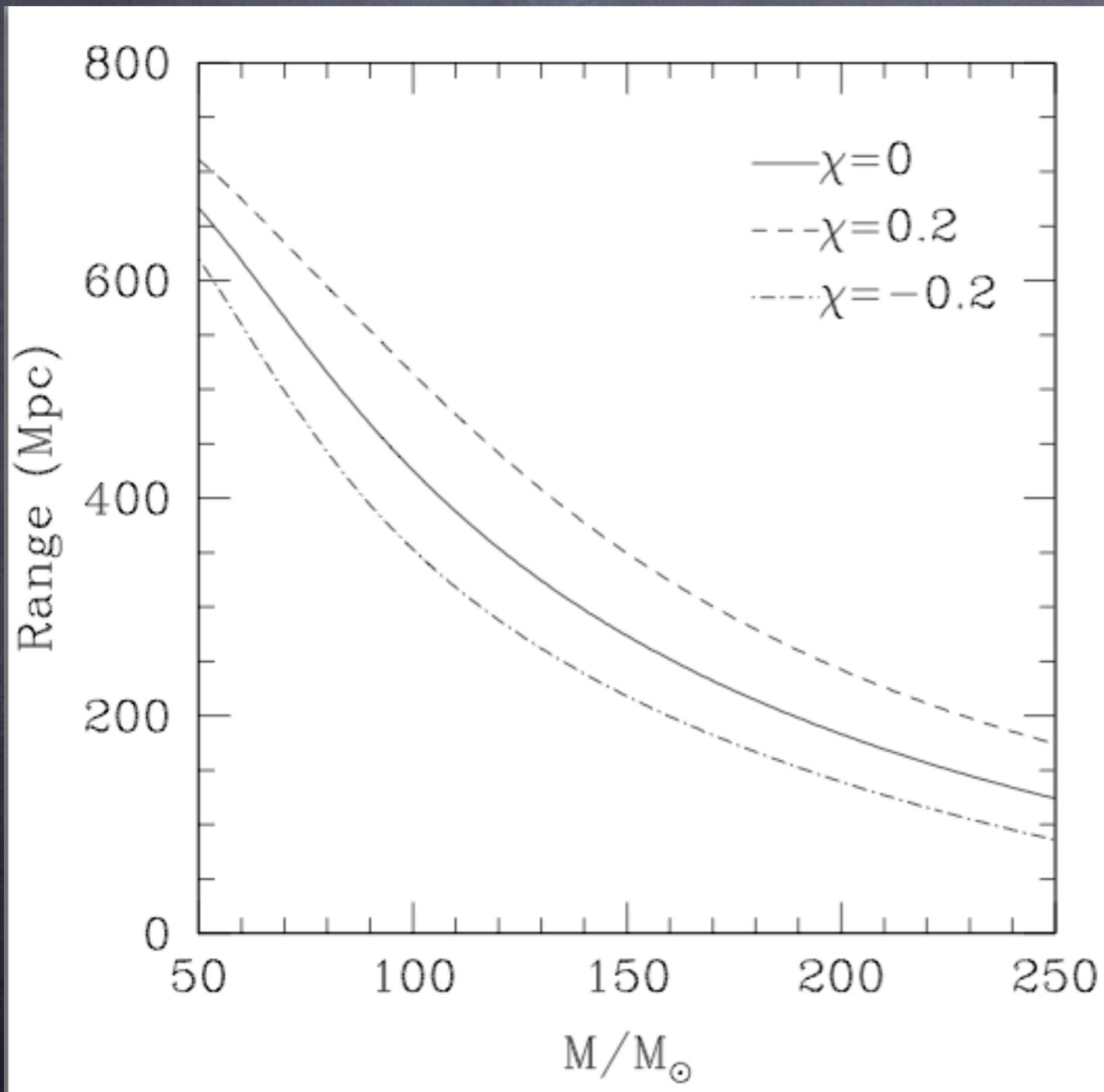
Event Rates: Mechanisms

- Three-body interactions: IMBH swaps into binaries, forms CO-IMBH binaries which are tightened via three-body interactions with other stars, then merge via GW radiation reaction
- Direct capture via energy loss to GWs
- Kozai resonances in hierarchical triple systems: inner binary eccentricity is driven up by outer companion
- Tidal capture of MS star that evolves into CO while in orbit
- Tidal interactions (orbital-vibrational coupling) for NS inspirals

Event rates per G.C.

- Binary tightening via 3-body interaction
- 3-body interaction rate is $dN/dt = n\sigma v$;
 $n \sim 10^{5.5} \text{ pc}^{-3}$; $v \sim 10 \text{ km/s}$; $\sigma \sim \pi a (2GM/v^2)$
- $T_{\text{harden}} \sim O(M/m) (dN/dt)^{-1} \sim 1.5 * 10^8 (AU/a) \text{ yr}$ [Quinlan]
- $T_{\text{merge}} \sim 5 * 10^{17} M_{\odot}^3 / (M^2 m) (a/AU)^4 (1-e^2)^{7/2} \text{ yr}$
 $\sim 5 * 10^8 (M_{\odot}/m) (100 M_{\odot}/M)^2 (a/AU)^4 \text{ yr}$ [Peters & Mathews]
- To maximize rate, minimize $T = T_{\text{harden}} + T_{\text{merge}}$
- Rate per globular is $\sim 3 * 10^{-9} \text{ yr}^{-1}$ for NS,
 $5 * 10^{-9} \text{ yr}^{-1}$ for BH

Advanced LIGO IMRI sensitivity



- Use EMRI-like waveforms, including non-quadrupolar harmonics, to determine range

- Range is spin-dependent

$$R \approx \left[1 + (\chi^2/2) \left(\frac{M}{100 M_{\odot}} \right)^{1.5} \right] \sqrt{\frac{m}{M_{\odot}}} \left[800 - 540 \left(\frac{M}{100 M_{\odot}} \right) + 107 \left(\frac{M}{100 M_{\odot}} \right)^2 \right]$$

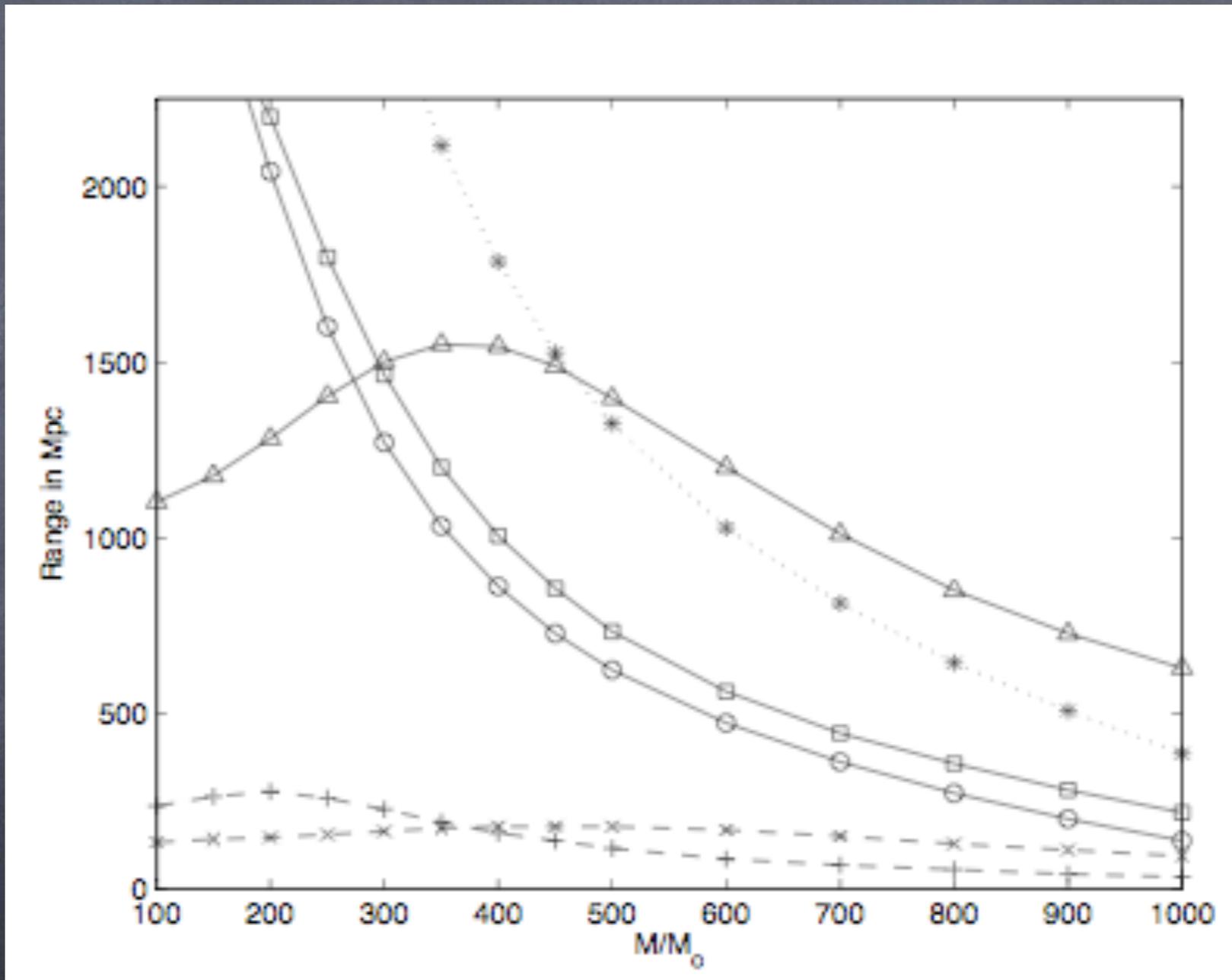
- Range could be increased by x1.5 by tuning Advanced LIGO

Advanced LIGO

IMRI rates

- Assume 10% of all globular clusters hold suitable IMBH (typical mass 100 Msun, spin=0.2)
- If inspiraling object is 1.4 Msun NS, Advanced LIGO could detect one IMRI per 3 years
- If inspiraling object is 10 Msun BH, Advanced LIGO could detect 10 IMRIs per year
- If Advanced LIGO is IMRI-optimized, rates could go up to 1/year and 30/year

Ringdowns



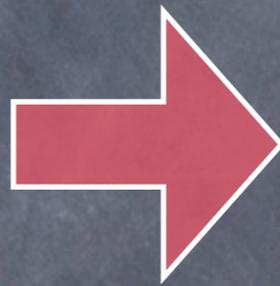
Could complement IMRIs if higher CO and IMBH masses are prevalent

Outline

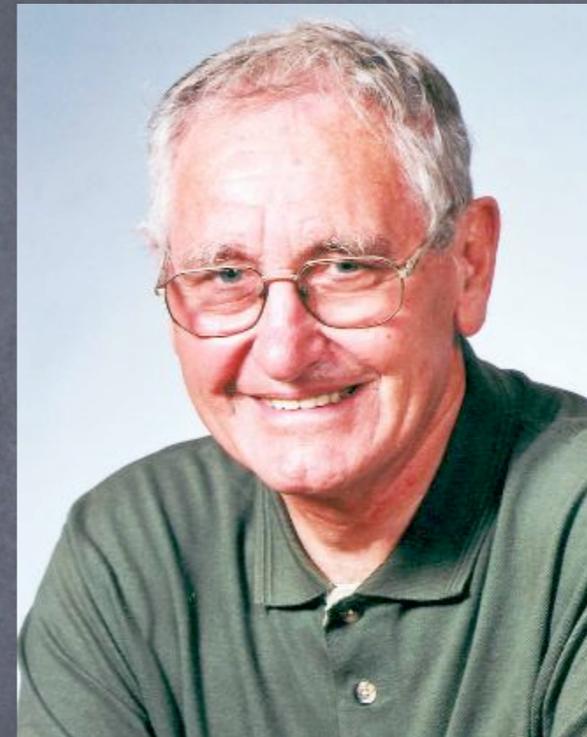
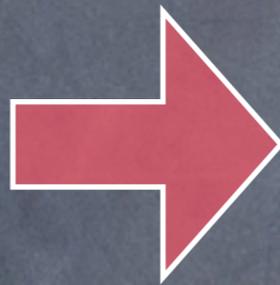
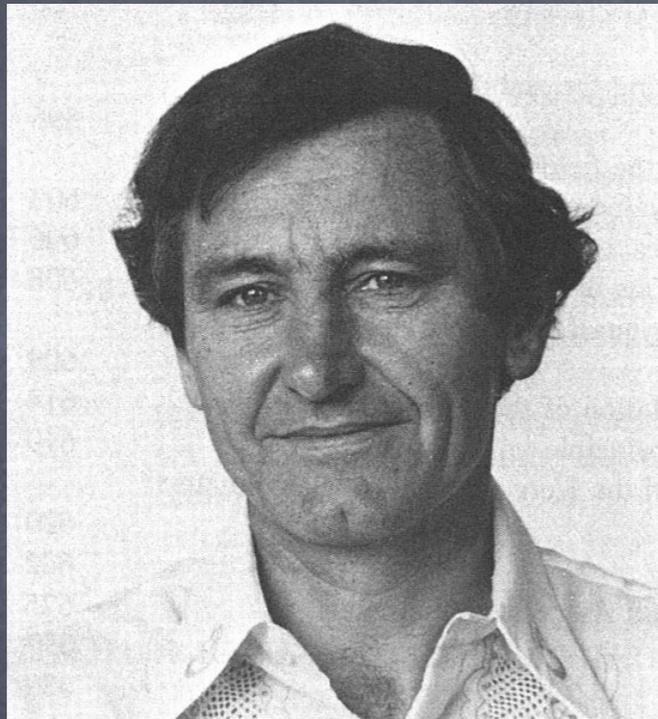
- GWs in ground-based detectors
- Intermediate-mass-ratio inspirals into intermediate-mass black holes: rates and characteristics
- Probing the strong field region near a black hole: testing the no-hair theorem with gravitational-wave observations

What is the "no-hair
theorem"?

What is the “no-hair theorem”?



What is the “no-hair theorem”?



Stationary, vacuum, asymptotically flat spacetimes in which the singularity is fully enclosed by a horizon with no closed timelike curves outside the horizon are described by the Kerr metric

The no-hair theorem in English

- “Black holes have no hair” means that all higher-order mass and current multipole moments are uniquely determined by the black hole mass and spin
- Conversely, an object with hair is one for which $M_n + iS_n \neq M(ia)^n$
- The “no-hair theorem” is a mathematical statement, so the title is a bit of a misnomer...

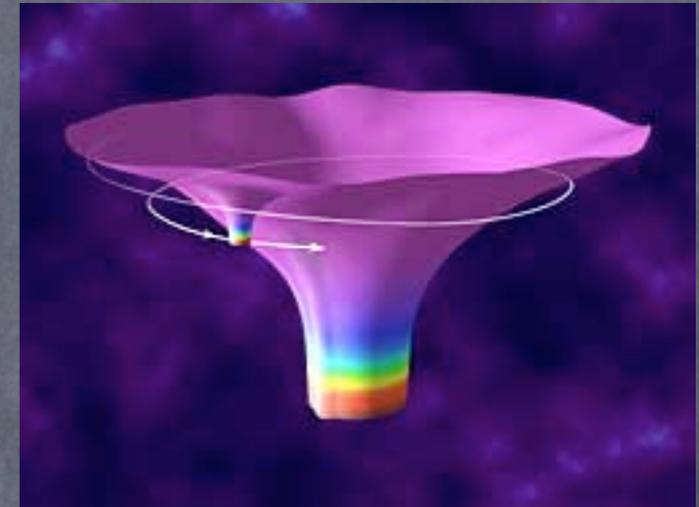
New subtitle: do massive black holes have hair?

- Are massive "black holes" really black holes?
- Could they be boson stars, or naked singularities, or...?
- Need to measure 3 multipole moments to test "Kerrness", 4 to test if an object is a boson star
- Search for exotic massive compact objects, test of cosmic censorship conjecture, null hypothesis test of the no-hair theorem...

Do Black Holes Have Hair?

Probing spacetime with E/IMRIs

- Are massive “black holes” really hairless?
- Or could they be boson stars, naked singularities, ...?
- Need to measure 3 multipole moments to test “Kerrness”, 4 to test if an object is a boson star
- LISA EMRIs into SMBHs will be the best probes of the strong-field regime ($\# \text{cycles} \sim M/m$), but Advanced LIGO IMRIs into IMBHs may provide the first interesting test
- Information about the spacetime structure and the orbit should be contained in GWs; how do we access it?



Summary

- Advanced LIGO could detect a few IMRIs per year
- Eccentricities will be low, circular waveforms can be used for detection (But should we use EMRI waveforms? Hybrid waveforms? ...?)
- Gravitational waves from EMRIs should make it possible to test whether the central body [SMBH] is a Kerr black hole
- Chaos in a non-Kerr spacetime would be an obvious smoking gun, but chaotic regions are probably not accessible
- Location of ISCO, periapsis precession, and orbital-plane precession are possible observables indicating bumpiness
- Frequency evolution over inspiral would be another observable, but more work is required

Event rates – upper limit

- Model-independent upper limit
- One core-collapsed globular cluster per Mpc^3
- One suitable IMBH per globular cluster
- IMBH grows from 50 to 350 solar masses by capture of COs in Hubble time
- Advanced LIGO could see IMRIs up to ~ 1000 Mpc (depending on masses, spin)
- Advanced LIGO may see tens of IMRIs per year (only 1 in 1000 years with Initial LIGO)
- Issues: kicks above 50 km/s eject IMBH; lower rates late in cluster history [e.g. simulations by O'Leary et al. 2006]

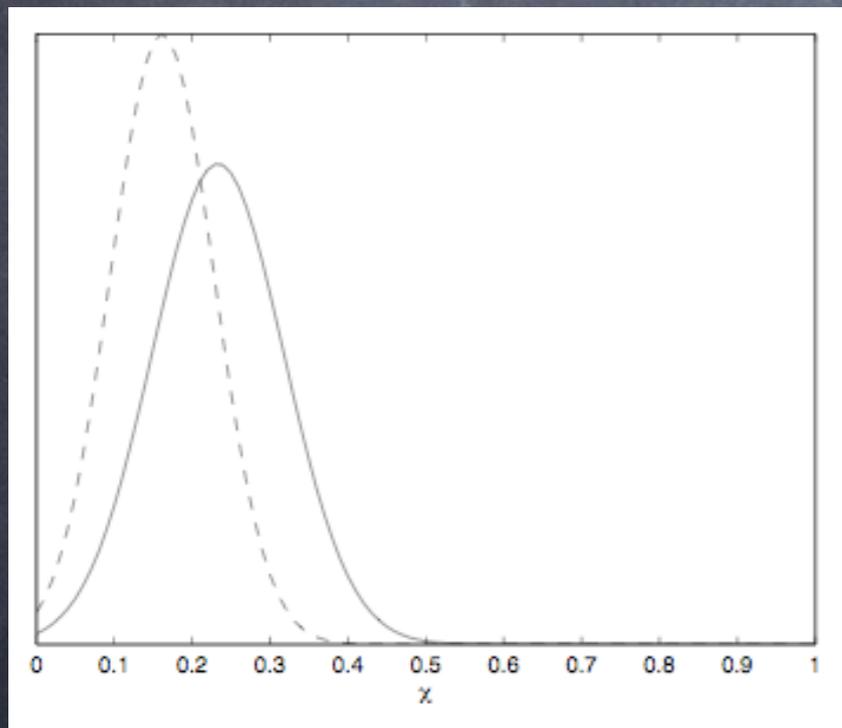
Spin and detection range

Evolution of spin distribution via minor mergers

Range increase due to spin

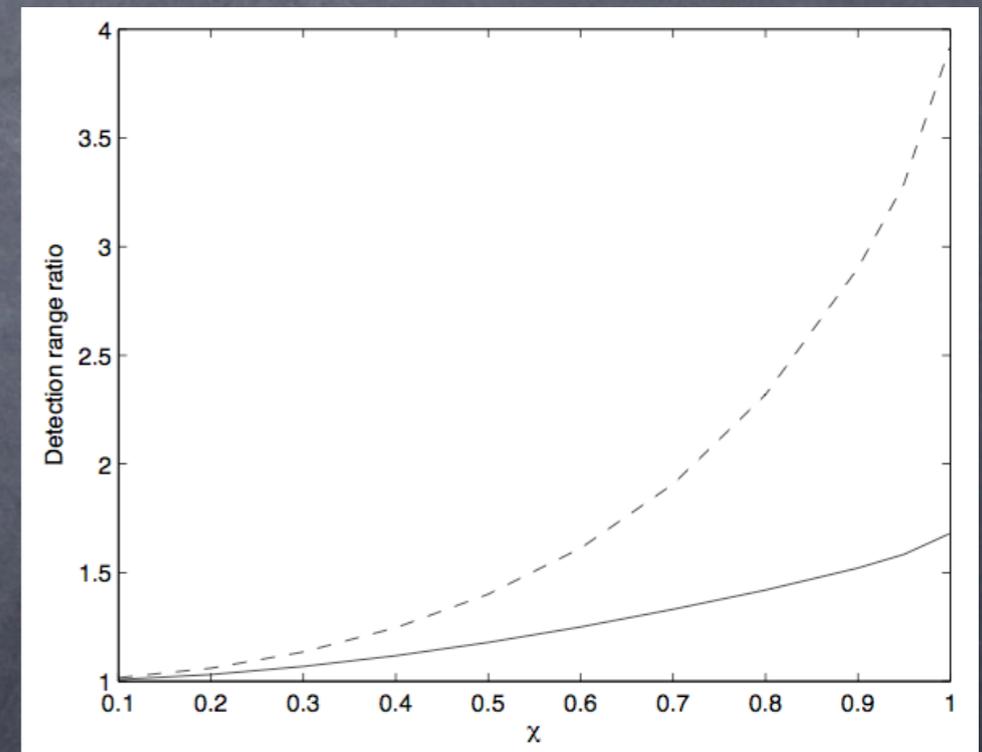
$$\frac{\partial}{\partial t} f(\chi, t) = -\frac{\partial}{\partial \chi} \left[\frac{\chi}{t} \left(-2 - \frac{4\sqrt{2}}{9} + \frac{4}{\chi^2 t} \right) f(\chi, t) \right] + \frac{1}{2} \frac{\partial^2}{\partial \chi^2} \left[\frac{4}{t^2} \left(1 + \frac{4\sqrt{2}\chi^2}{9} - \chi^2 \right) f(\chi, t) \right]$$

$$\frac{\text{Range}_{\text{spin}}}{\text{Range}_{\text{no-spin}}} \sim 1 + 0.6\chi^2 \left(\frac{M}{100 M_{\odot}} \right)$$



$t=M/m$

$$\bar{\chi} \approx \bar{\chi}_0 \left(\frac{t}{t_0} \right)^a \approx \bar{\chi}_0 \left(\frac{M_0}{M} \right)^{2.63}$$



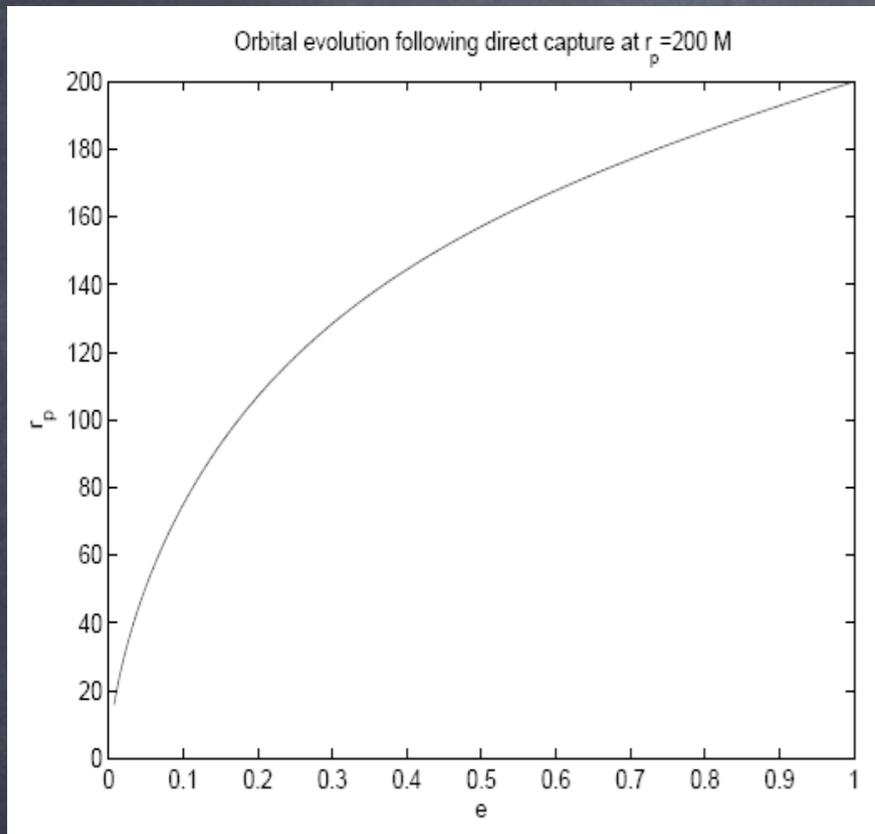
Evolution from $t=M/m=50$ to $t=100$ (e.g., from $M=70$ to $M=140$ solar masses via capture of $m=1.4$ solar-mass NSs)

If initial $\chi=0.1$, then mean spin at $t=100$ is 0.162, $\sigma=0.066$

If initial $\chi=0.9$, then mean spin at $t=100$ is 0.233, $\sigma=0.087$

Solid line - inspiral into 100 Msun IMBH
Dashed line - inspiral into 200 Msun IMBH
Effect is very pronounced for LISA:
can cause bias in spin estimate

Eccentricities in AdvLIGO band



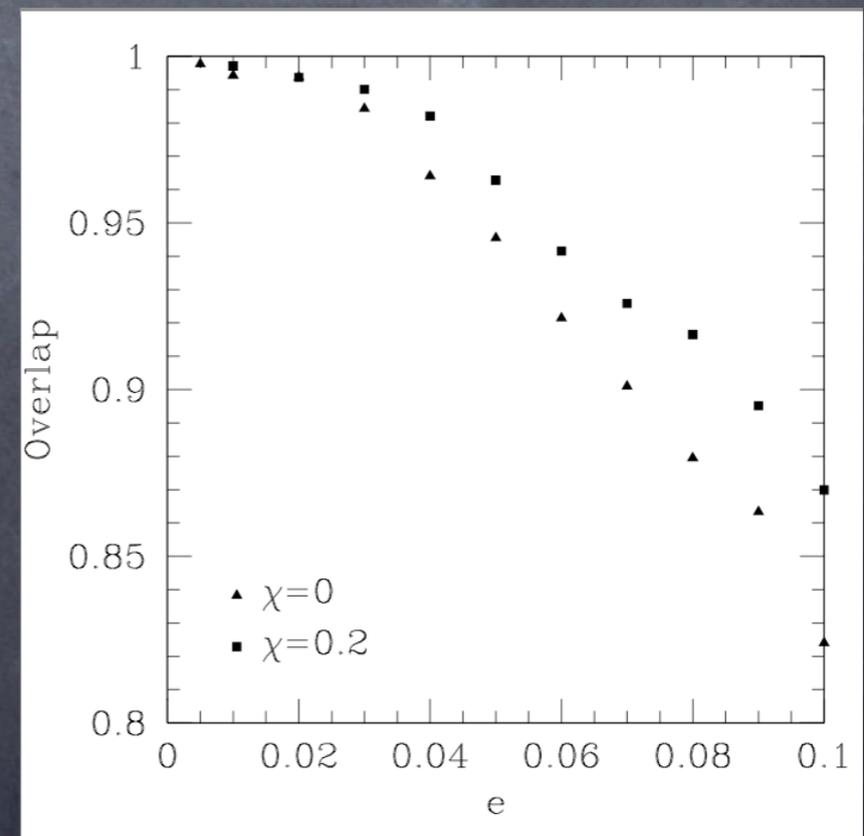
- Hardening via 3-body interactions

Eccentricity $\sim \text{few} \times 10^{-5}$ when $f_{\text{GW}} = 10$ Hz

- Direct capture

90% of IMRIs circularize to $e < 0.1$ by 10 Hz, 67% circularize to $e < 0.01$ by $f_{\text{GW}} = 10$ Hz

- At $e = 0.01$, overlap between eccentric and circular templates is > 0.99 , so circular templates can be used for detection



Observing deviations from Kerr with EMRIs

- LISA can detect tens to thousands of EMRIs
- Ryan's theorem [1995]: GWs from nearly circular, nearly equatorial orbits in stationary, axisymmetric spacetimes encode all of the spacetime multipole moments... in principle
- Can we extend this theorem? Are there obvious observable imprints of an anomalous, non-Kerr quadrupole moment (a "bumpy" spacetime)?
- Are energy E , angular momentum L_z and Carter constant Q conserved in a bumpy spacetime?

Geodesics in bumpy spacetimes

- Use Manko–Novikov bumpy spacetime

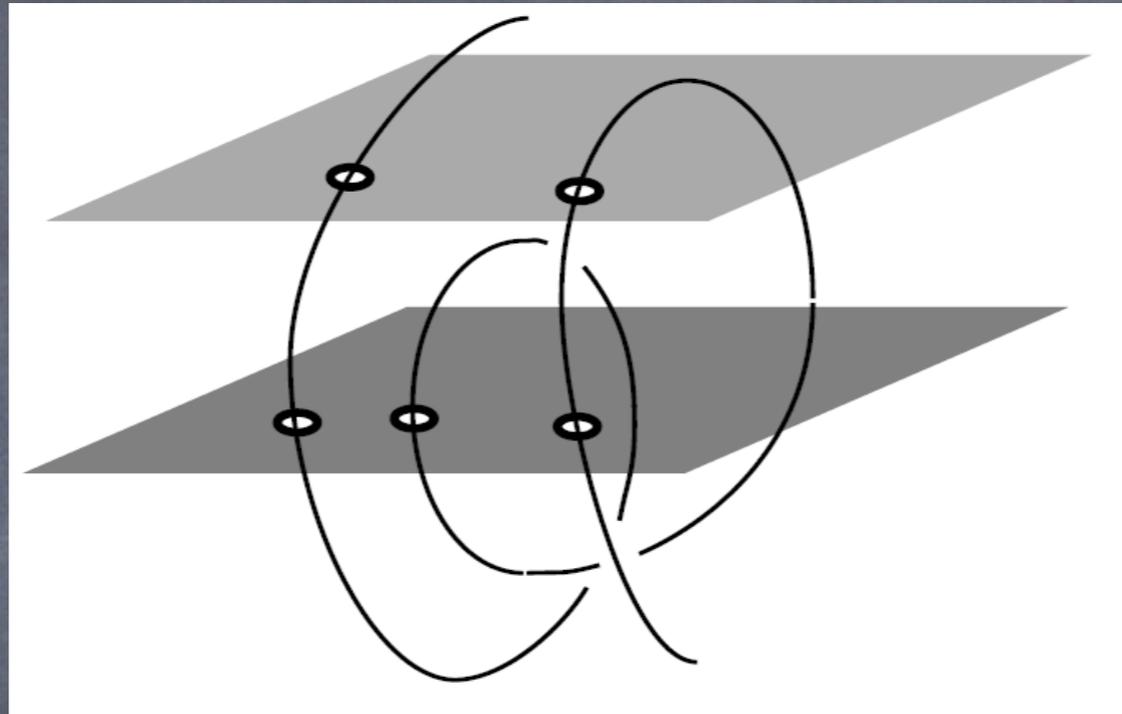
$$ds^2 = -f(\rho, z) (dt - \omega(\rho, z) d\phi)^2 + \frac{1}{f(\rho, z)} \left[e^{2\gamma(\rho, z)} (d\rho^2 + dz^2) + \rho^2 d\phi^2 \right]$$

- C code - geodesic equations:

$$\frac{\partial^2 x^\alpha}{\partial \tau^2} = -\Gamma_{\beta\gamma}^\alpha \frac{\partial x^\beta}{\partial \tau} \frac{\partial x^\gamma}{\partial \tau}$$

- Check conservation of E , L_z , 4-velocity norm
- Equations might not separate as in Kerr
- Is there a full set of integrals of motion?

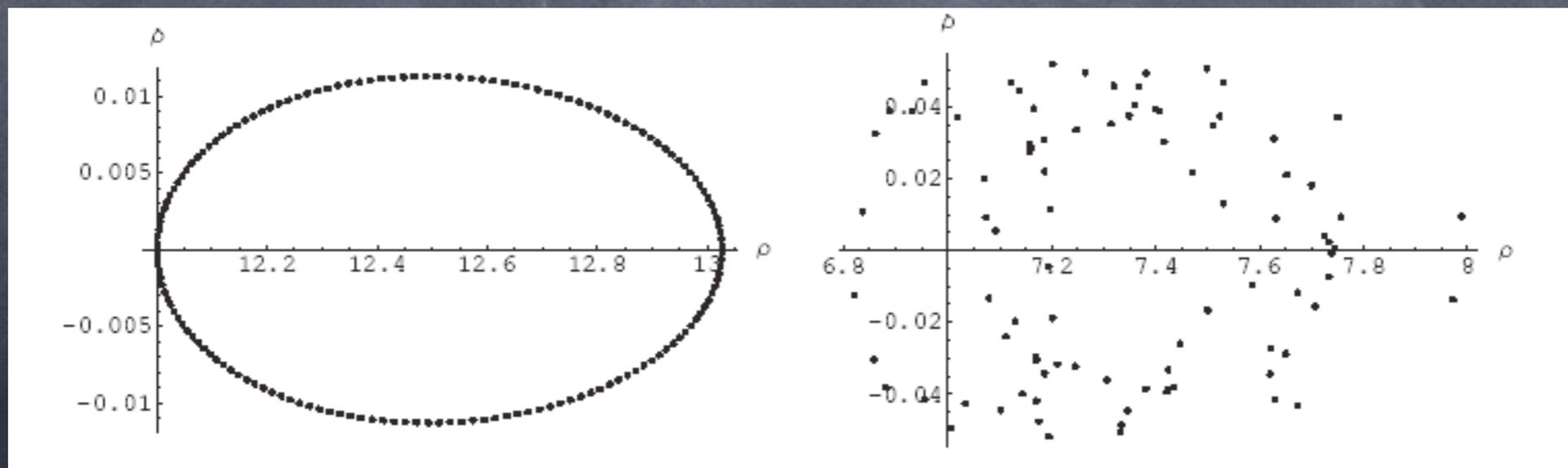
Poincare maps



- Check if spacetime has a full set of integrals of motion
- Plot dp/dt vs. p for $z=z_0$ crossings
- Phase space plots should be closed curves for all z_0 iff there is a third isolating integral

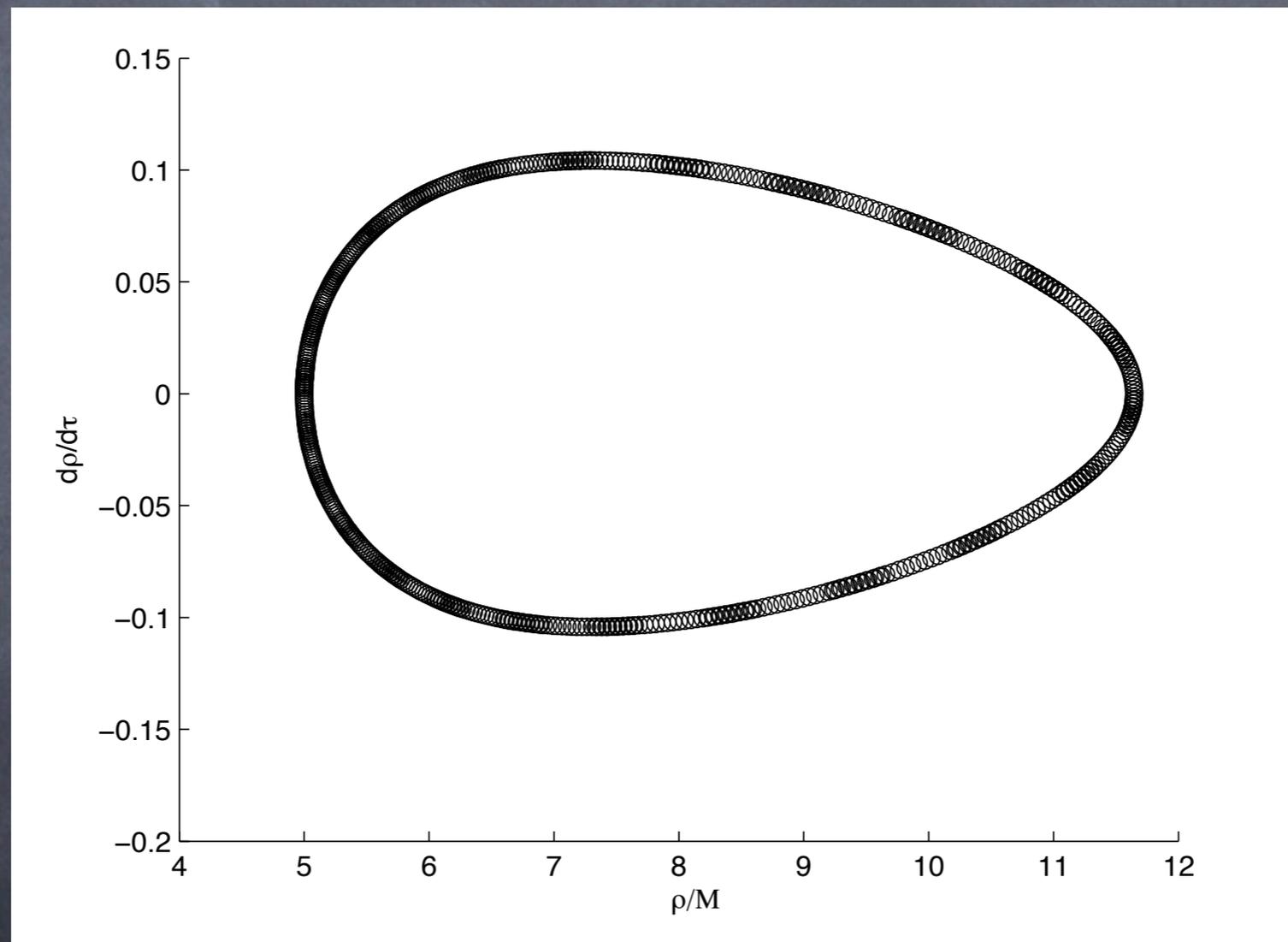
Poincare maps for motion in Newtonian potential with hexadecapole moment

$$V(r, \theta) = -\frac{M_0}{r} + \frac{M_2}{r^3} P_2(\cos \theta) + \frac{M_4}{r^5} P_4(\cos \theta)$$



$$M_2 = 10 M_0; \quad M_4 = 400 M_0$$

Poincare map in a bumpy spacetime



$$E=0.95, L_z=-3, a/M=0.9, q=0.95$$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$

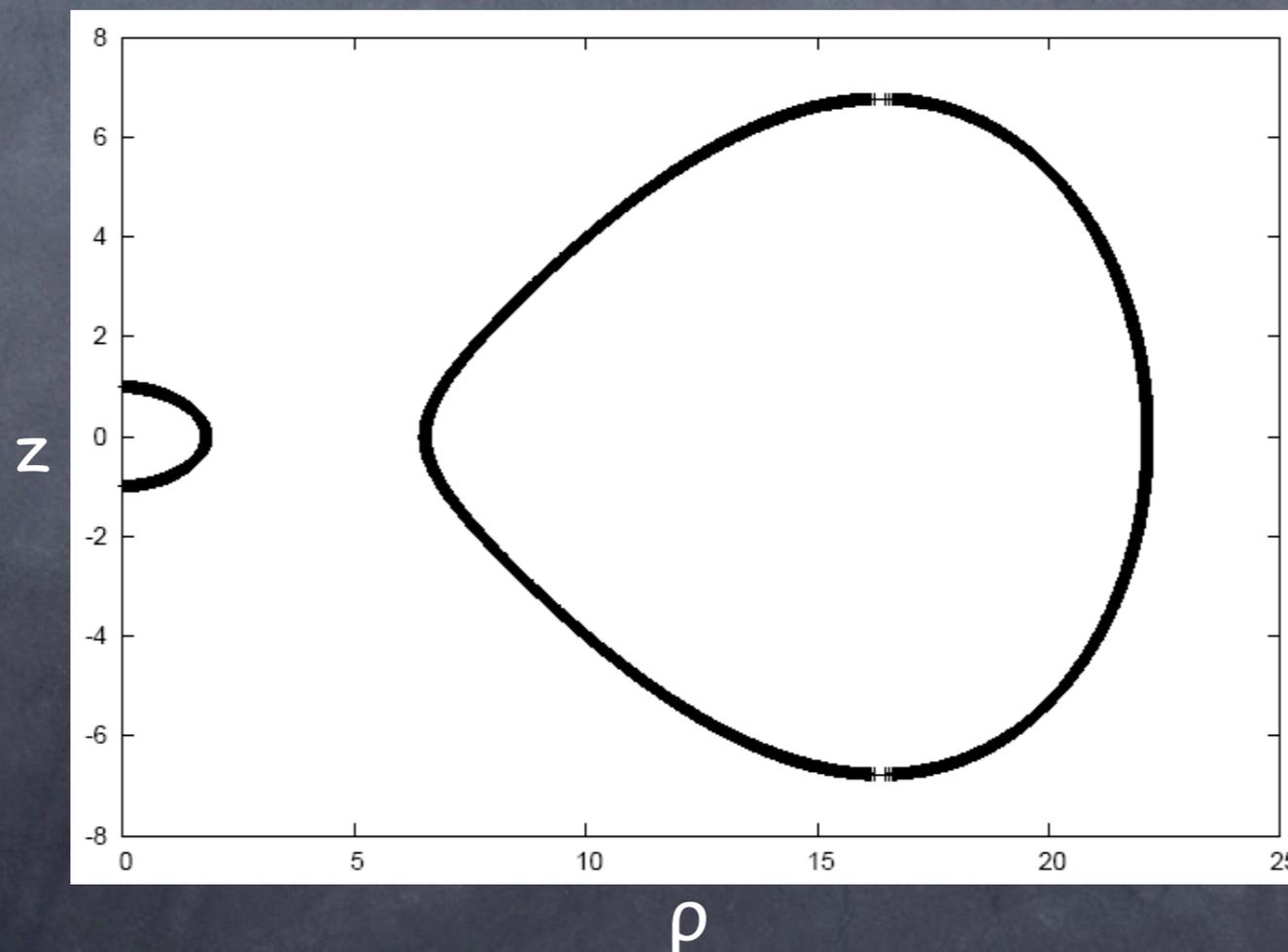
z

ρ

$E=0.95, L_z=-3, a/M=0.9$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$



$E=0.95, L_z=-3, a/M=0.9$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$

z

ρ

$E=0.95, L_z=-3, a/M=0.9$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$

z

z

ρ

ρ

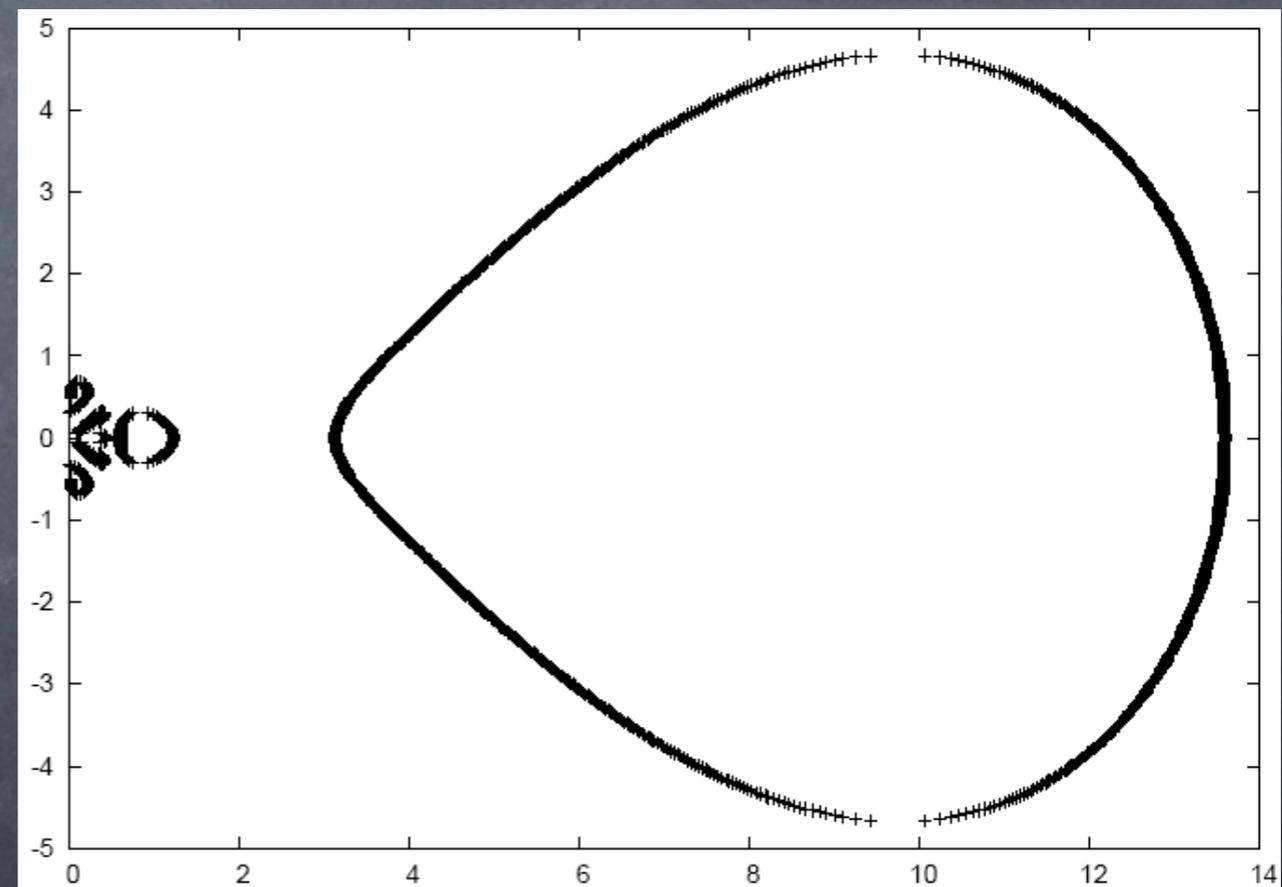
$E=0.95, L_z=-3, a/M=0.9, q=0.95$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$

z

z



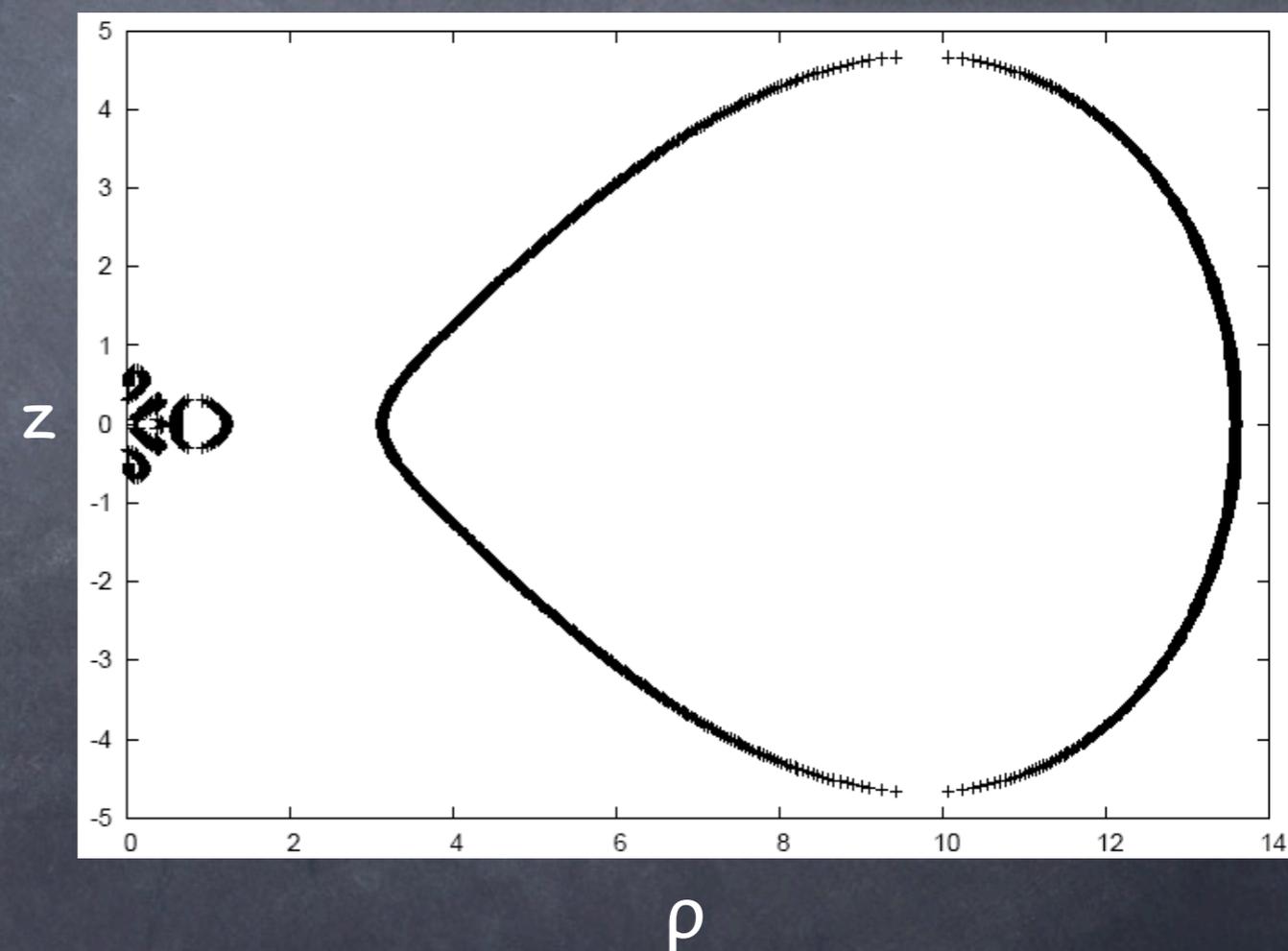
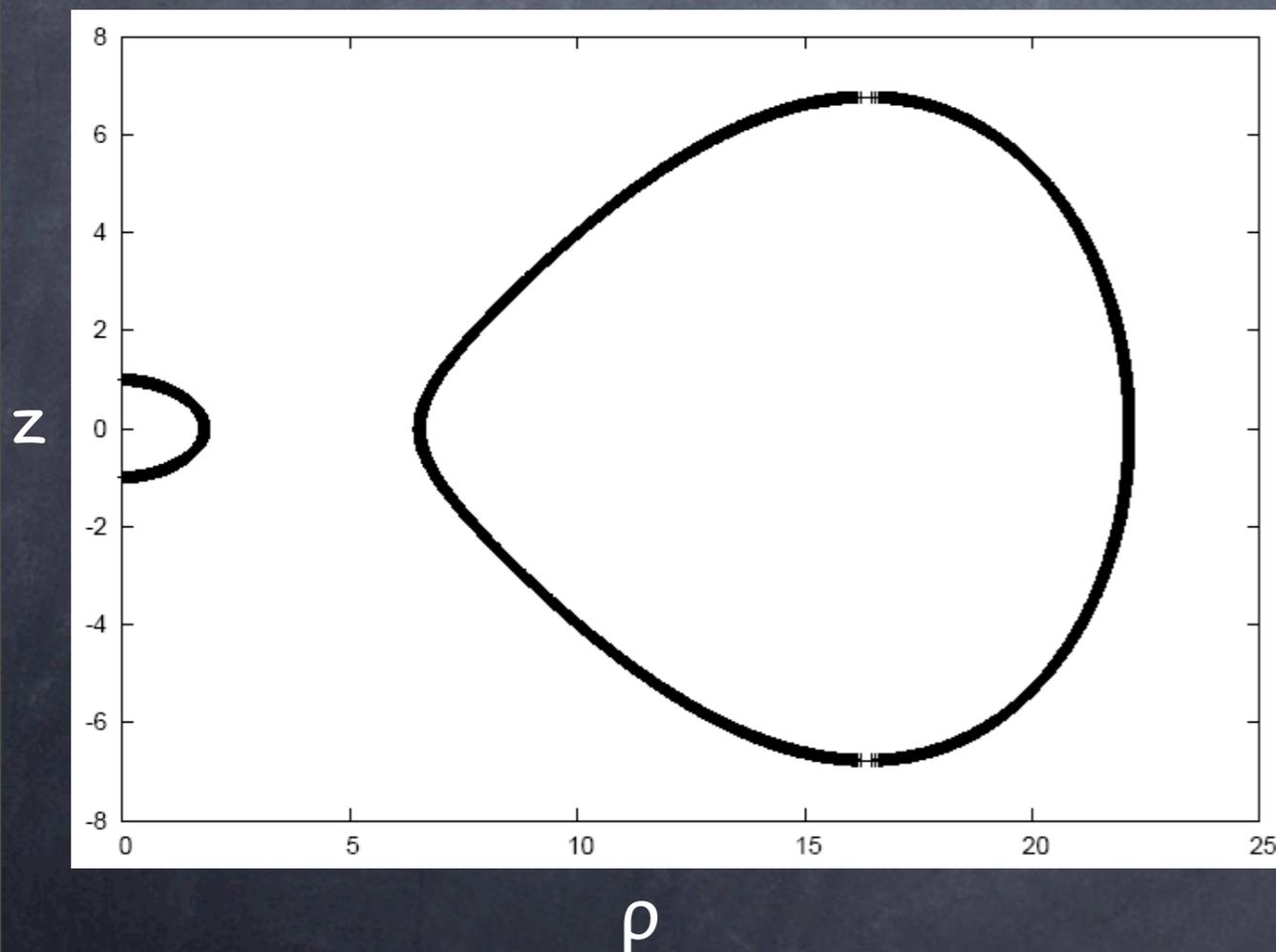
ρ

ρ

$E=0.95, L_z=-3, a/M=0.9, q=0.95$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$



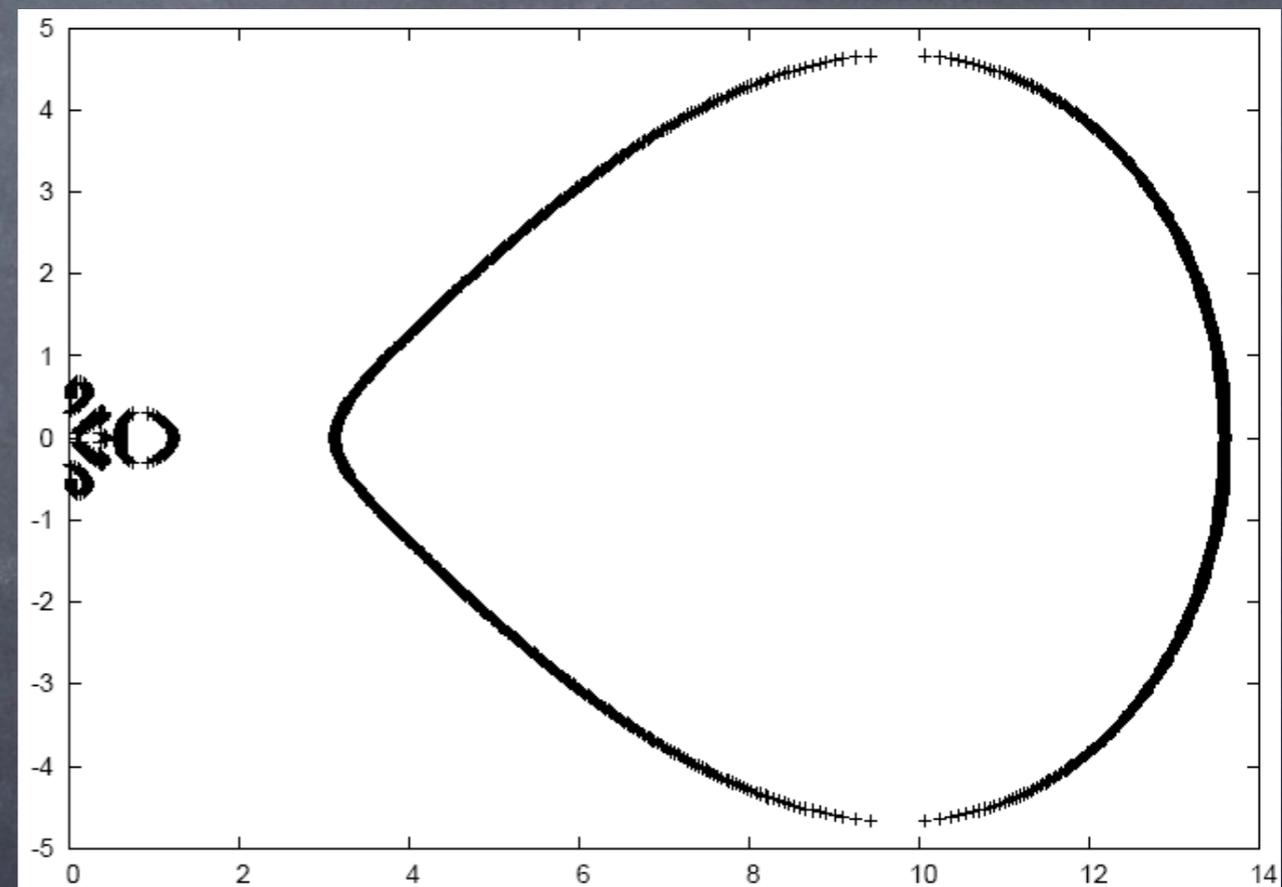
$E=0.95, L_z=-3, a/M=0.9, q=0.95$

Allowed regions for bound orbits

Effective potential $(\dot{\rho}^2 + \dot{z}^2) = V(E, L_z, \rho, z)$

z

z

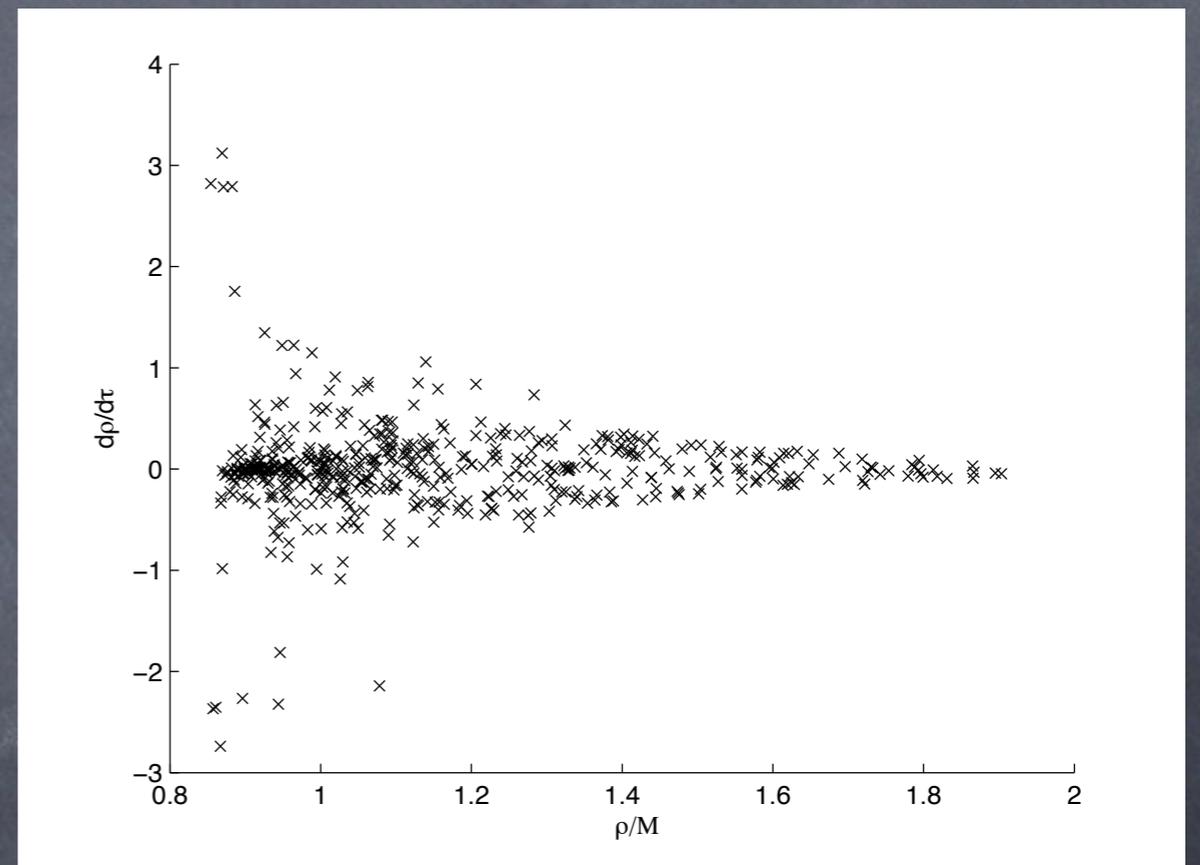
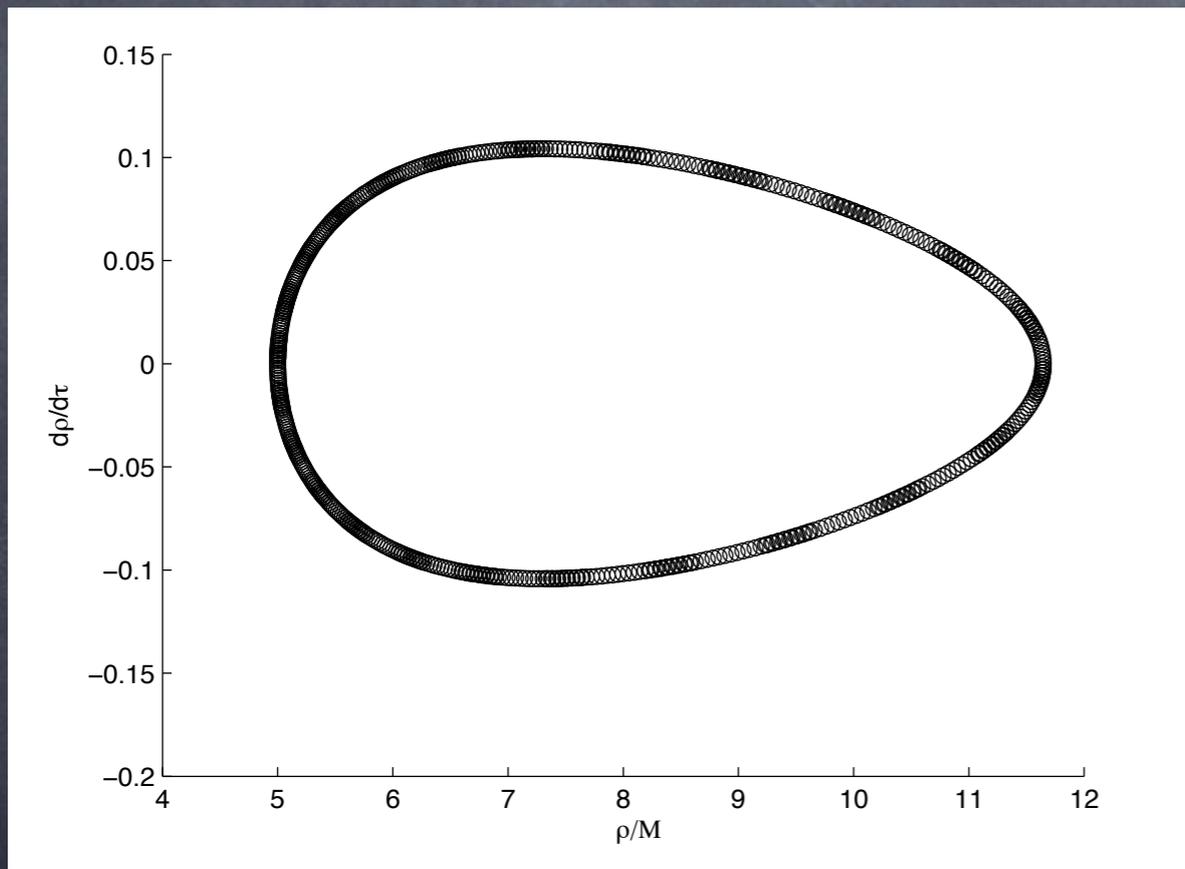


ρ

ρ

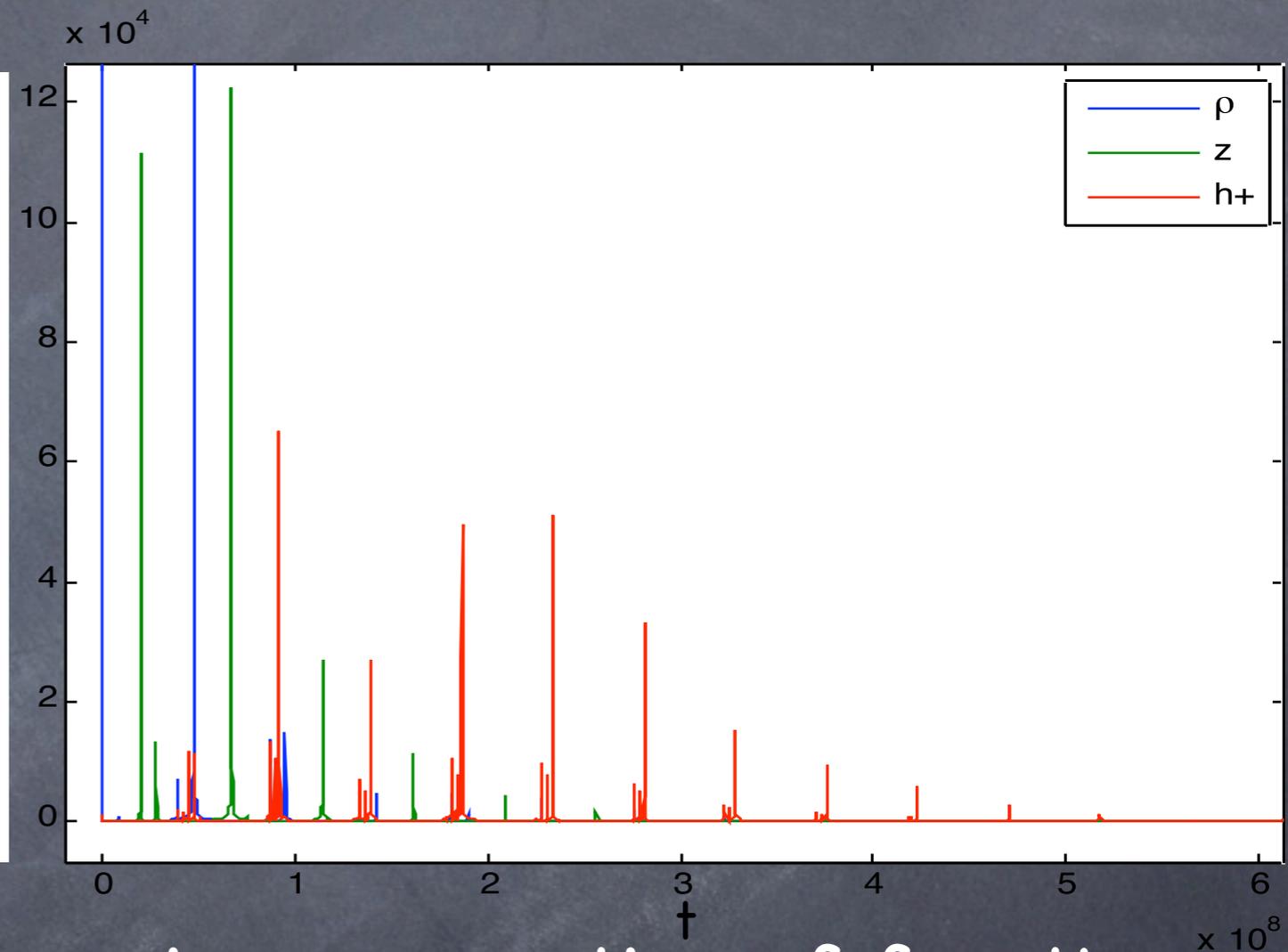
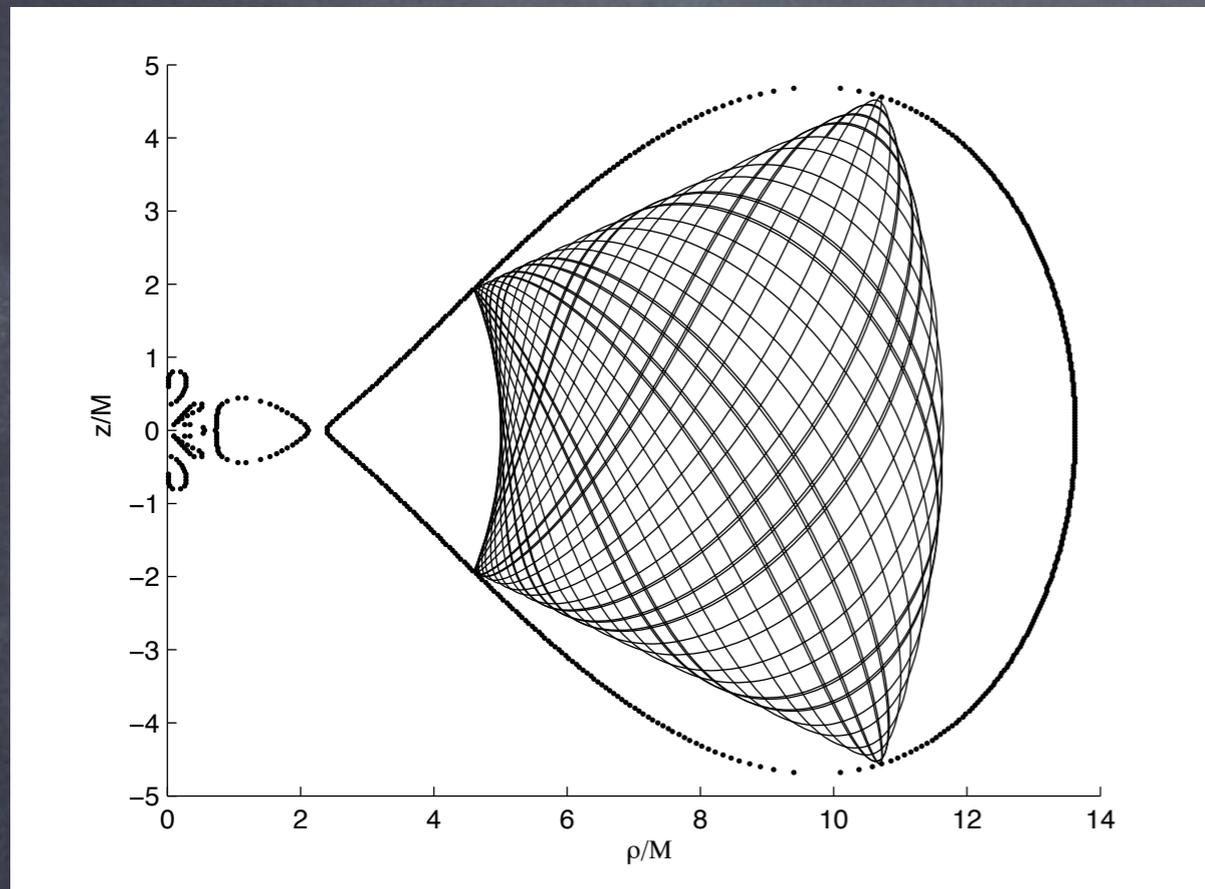
$E=0.95, L_z=-3, a/M=0.9, q=0.95$

Poincare map in a bumpy spacetime, second look



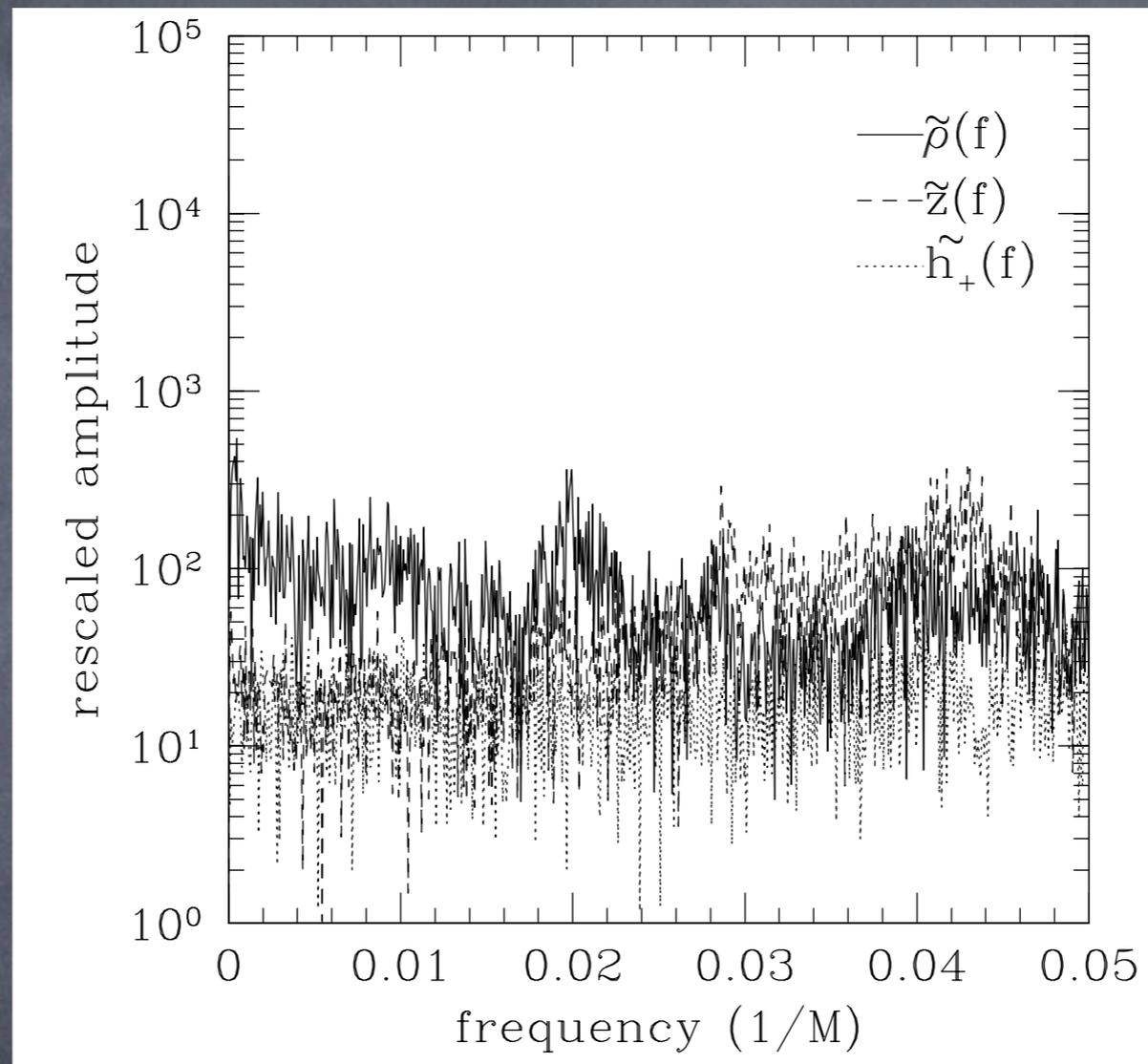
$E=0.95, L_z=-3, a/M=0.9, q=0.95$

Regular outer region



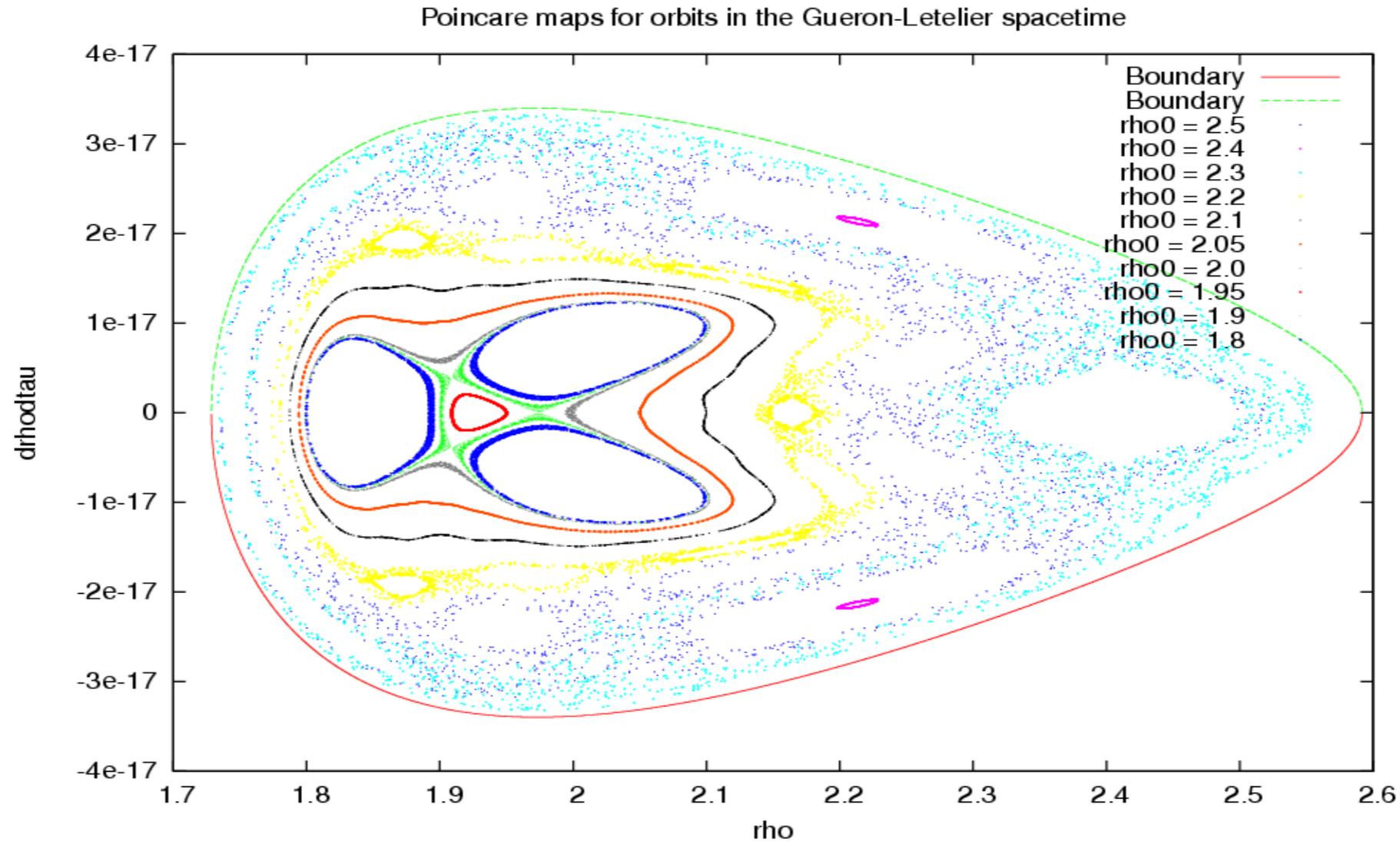
- Regular motion in outer region, suggestive of fourth-degree invariant
- Both ρ and z motion consist of harmonics of two fundamental frequencies to 10^{-7}

Chaotic inner region

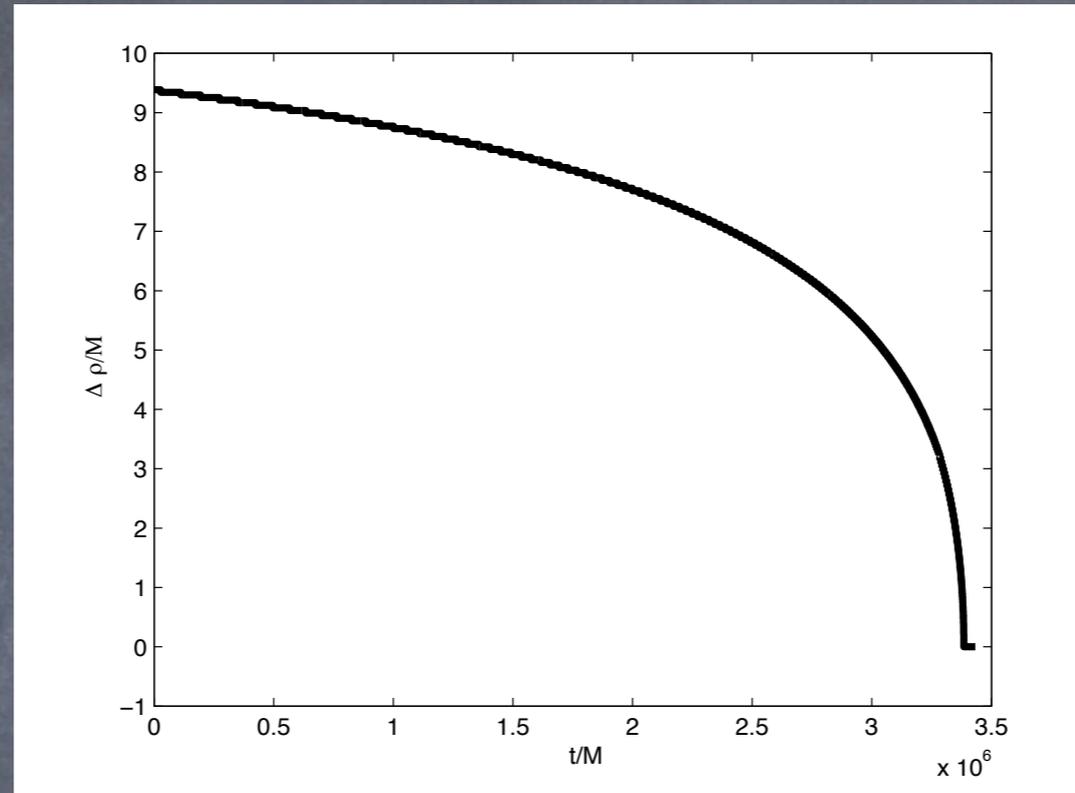


- If motion is chaotic for any initial conditions, it is chaotic for all initial conditions, but an approximate invariant may exist in some cases (invariant tori) [KAM Theorem]

Chaos in Gueron-Letelier spacetime



Is chaos accessible?

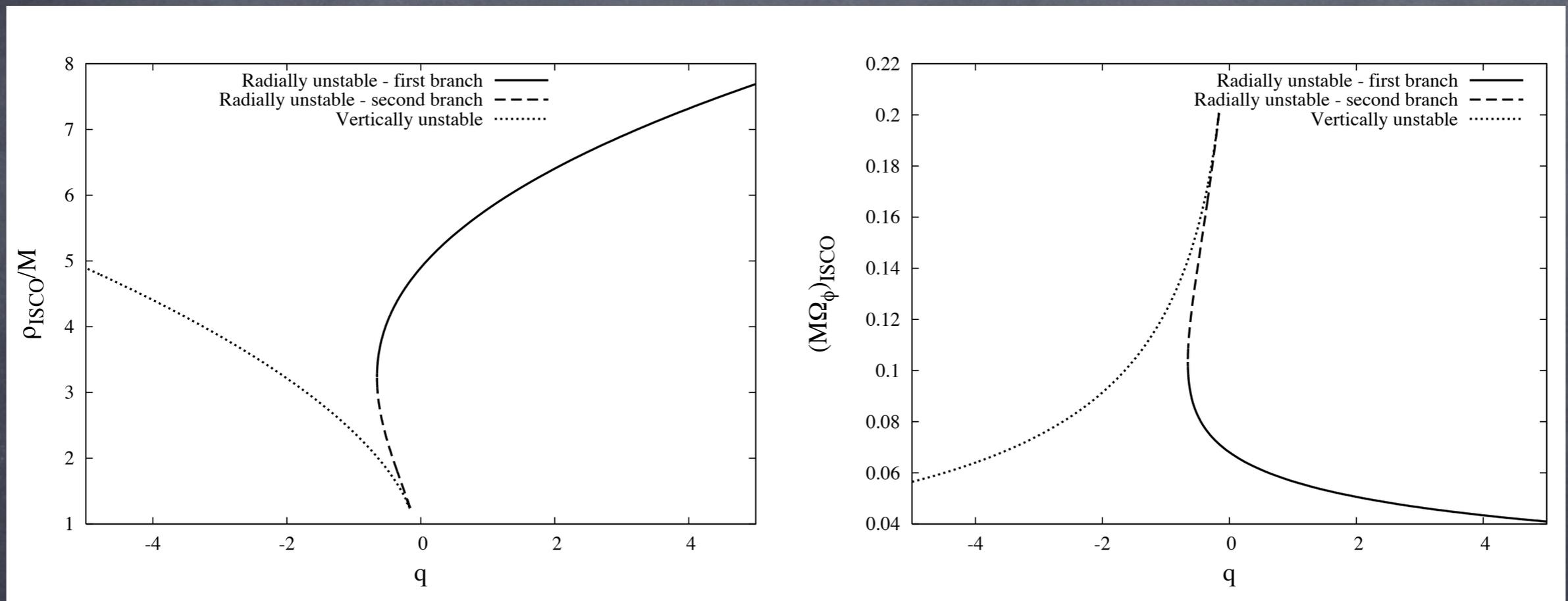


- Inner and outer regions appear to merge under radiation reaction, but never split
- Object starts out in outer, regular region; once the two regions are fully merged, motion is regular (but odd things may happen when the neck is narrow...)

Other observable signatures of bumpiness

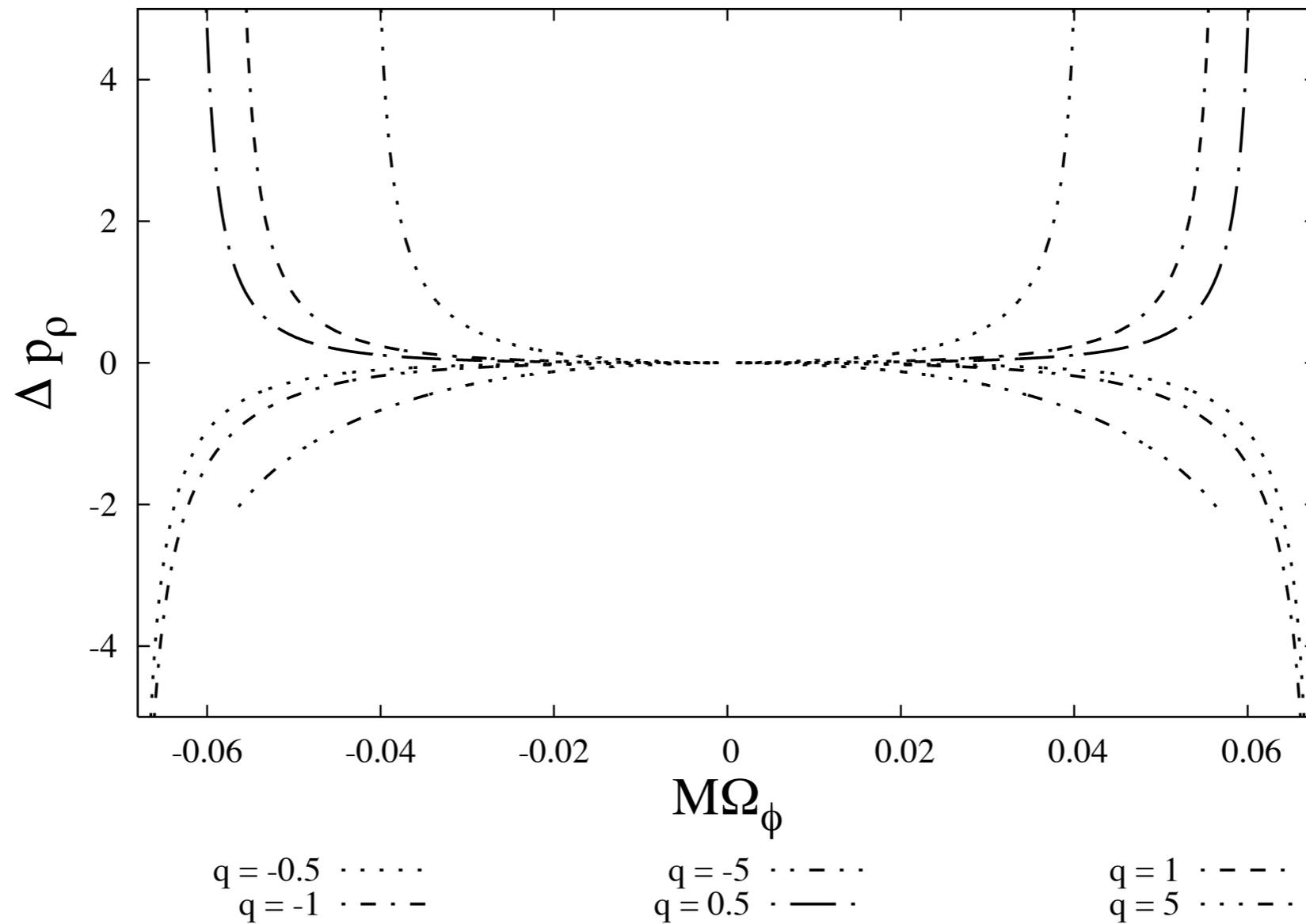
- If the orbits are indeed multi-periodic, then the spacetime “bumpiness” should be observable via:
 1. three fundamental frequencies of gravitational waves
 2. harmonic structure of the waves (relative frequencies and phases of harmonics)
 3. evolution of these with time over inspiral
- Further study required to properly analyze inspiral

Location of innermost stable circular orbit (ISCO)

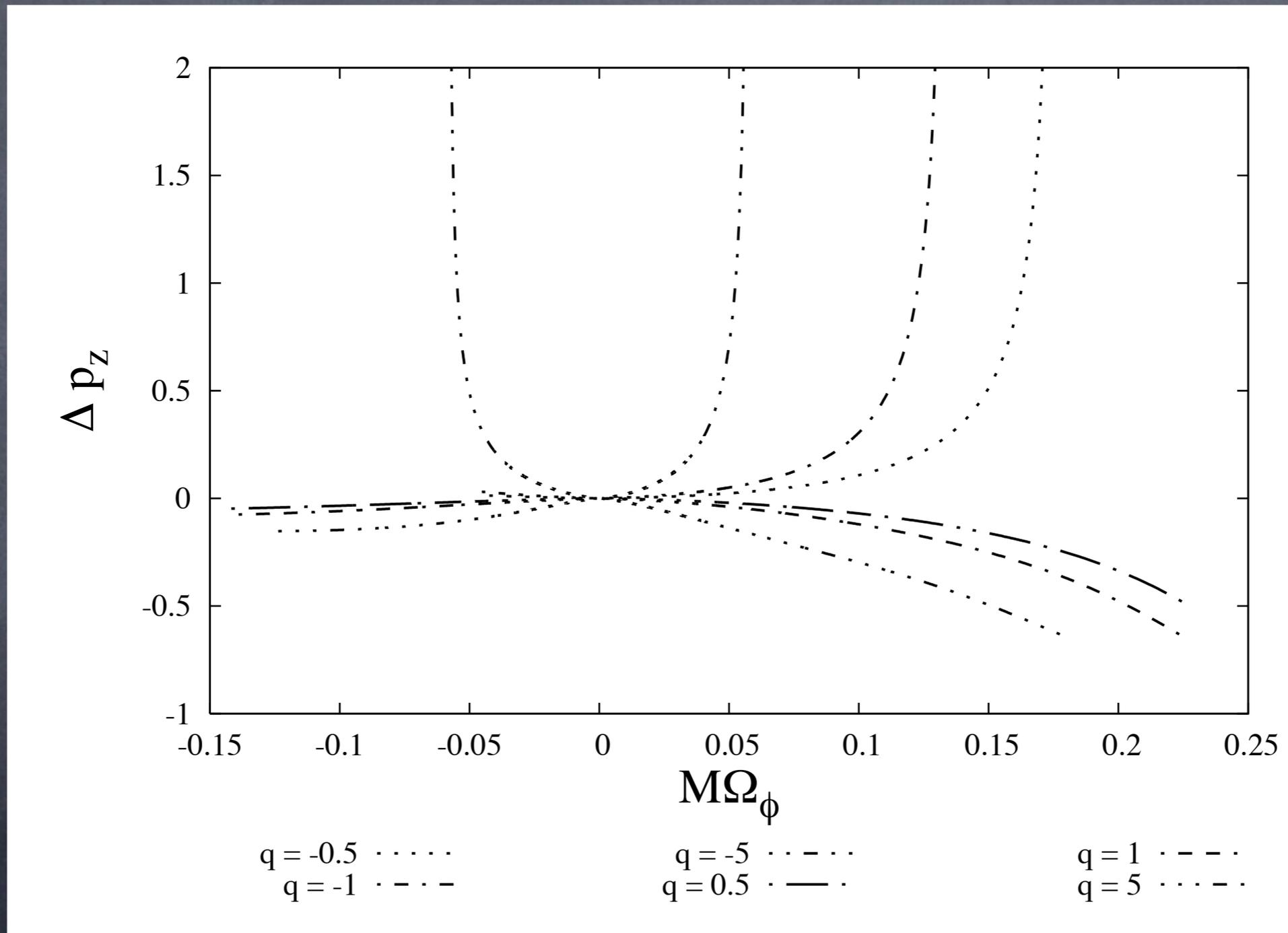


- The ISCO frequency (and hence plunge frequency) depends on the value of the spacetime quadrupole moment

Periapsis precession



Orbital-plane precession



Summary

- Advanced LIGO could detect a few IMRIs per year
- Eccentricities will be low, circular waveforms can be used for detection (But should we use EMRI waveforms? Hybrid waveforms? ...?)
- Gravitational waves from EMRIs should make it possible to test whether the central body [SMBH] is a Kerr black hole
- Chaos in a non-Kerr spacetime would be an obvious smoking gun, but chaotic regions are probably not accessible
- Location of ISCO, periapsis precession, and orbital-plane precession are possible observables indicating bumpiness
- Frequency evolution over inspiral would be another observable, but more work is required

Do I really believe that IMBHs exist and MBHs are not black holes?

- I don't know. But it's dangerous to assume that one will see only what one expects to see. We should be prepared to test our assumptions.
- Every time a new part of the electromagnetic spectrum was accessed (radio-astronomy, X-rays, etc.), something unexpected was seen. Gravitational waves are a new window to the universe: expect to see the unexpected!